

Research Article Maxwell's Quaternion Equations

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Abstract

The equations of electrodynamics must, first of all, satisfy the law of conservation of energy. It is shown that Maxwell's equations can be obtained from the Cauchy-Riemann conditions for a quaternion in 4D space. Electrons are written as 4D vectors in energy space, in which the first elements represent the real part of the quaternion (scalar), and the other three represent the imaginary part. From the point of view of conservation of energy, an electron cannot move into an arbitrary state, but only makes quantum jumps to those places in space in which it stores energy. Consequently, the movement of electrons in time occurs along an orbit. Since scalars are formed by the interaction of electromagnetic waves, 4D electrons have a spectrum.

The mathematically obtained equations of quaternion electrodynamics have the same form for electric and magnetic intensity, but differ from Maxwell's equations by the presence of a scalar part. A charged electron is considered as the scalar part in the equation of circulation of the electric field strength. The electron spin is considered as the scalar part in the equation of magnetic intensity circulation. The equations of the scalar parts correspond to Gauss's law and form a single connection with the equations of the imaginary parts. Also, unlike Maxwell's equations, instead of currents induced by circulations of intensities, the electromotive forces that form these currents are shown. As is known, in the equation for the circulation of magnetic intensity, Maxwell added a current formed by the change in electric flux over time. In the obtained expressions, this term appeared mathematically and represents the electromotive force generated by the change in the magnetic field over time.

Keywords: Maxwell's Equations, Quaternion, Electrodynamics, Circulation, Rotor

1. Introduction

Electrodynamics is based on Maxwell's equations [1,2]. Maxwell generalized the experiments on electricity of Faraday and Ampere and described them in terms of vector mathematics in 1865. The discovery of the quaternion by William Rowan Hamilton occurred in 1843. In formulating the equations of electrodynamics, Maxwell used the Hamiltonian operator in 3D space with imaginary units *i*, *j*, *k*. This operator is a pure quaternion, i.e., a quaternion without a scalar part. The absence of the scalar part resulted in a zero solution to the equation based on Ampere's law. Therefore, Maxwell introduced into this equation the time derivative of the electric flux.

In everyday life, physical space is perceived as three-dimensional: length, width, height. In physics, time is added to the dimension of space. However, physical space is universal and allows for the formation of spaces of much greater dimensions within it. This fact is confirmed by the evolution of Nature on Earth and in outer space in general. The more complex the structure of matter, the greater the dimensionality of the space in which it is formed must be.

Mathematically, these phenomena can be described using hypercomplex numbers [3]. Quaternion forms a 4D space in 3D physical space, octonion -8D, sedenion -16D, etc. The dimension of the space is doubled by the dimension of the hypercomplex number. Doubling of dimensionality is associated with the procedure of doubling complex numbers when obtaining hypercomplex signals. Imaginary numbers *i*, *j*, *k*, etc. form orthogonal coordinate axes of such a space. The presence of scalar and imaginary parts in hypercomplex numbers allows us to explain how particles emerge from waves. Electromagnetic waves become material rather than a special kind of matter. The interaction of waves occurs not due to the ether, but directly due to the particles formed by them and vice versa.

The purpose of the article is to present a quaternion mathematical model of charges and their interaction with electric and magnetic fields, on the basis of which to obtain the equations of electrodynamics.

2. Materials and Methods for Solving the Problem

To derive the equations of electrodynamics, we use a quaternion. In algebraic form, a quaternion is written as [3,4]

$$q = s + ix + jy + kz,$$

where s, x, y, z are real numbers, i, j, k are imaginary units.

A quaternion function is also a quaternion, which can be written as a sum of functions p(s,x,y,z), u(s,x,y,z), v(s,x,y,z), w(s,x,y,z):

$$f(q) = p + iu + jv + kw.$$

Quaternion (1) and function of quaternion (2) have dimension 4D in 4D space with real coordinate axis s and three imaginary axes i, j, k. To eliminate imaginary units in mathematical models of systems in multidimensional space, it is proposed to use their vector and matrix representation [5].

In vector representation, the function of quaternion (2) will look like: $f(q) = \begin{bmatrix} p & u & v & w \end{bmatrix}^T$. In matrix representation, we write quaternion (1) as

$$\mathbf{Q} = \begin{bmatrix} s & x & y & z \\ -x & s & -z & y \\ -y & z & s & -x \\ -z & -y & x & s \end{bmatrix}.$$
 (3)

Let us extract the basis matrices from (3):

	1	0	0	0	[()	1	0	0		0	0	1	0		0	0	0	1]	
	0	1	0	0	- -	1	0	0	0	Ŧ	0	0	0	1	V	0	0	-1	0	
F =	0	0	1	0,	l = ()	0	0	-1	, J =	-1	0	0	0	, K =	0	1	0	0	. (4)
	0	0	0	1)	0	1	0		0	-1	0	0_		1	0	0	0	

Using the basis matrices (4), we represent the quaternion (1) as follows:

$$\mathbf{Q} = \mathbf{E}s + \mathbf{I}x + \mathbf{J}y + \mathbf{K}z$$

For quaternion (5) the multiplication table of basis matrices (4) shown in Table 1 is valid.

Table 1: Multiplication operations of quaternion basis matrices.

×	E	Ι	J	K
E	E	Ι	J	K
I	Ι	$-\mathbf{E}$	K	$-\mathbf{J}$.
J	J	$-\mathbf{K}$	$-\mathbf{E}$	Ι
K	K	J	$-\mathbf{I}$	$-\mathbf{E}$

As can be seen from the basis matrices (4), representation (5) and Table 1, the imaginary units and, accordingly, the axes of spatial coordinates, also have a dimension of 4D. The basis matrices are orthogonal and do not intersect in space, i.e., when they are superimposed on each other.

Let us consider the use of quaternion (3) as a differential operator for the equation of dynamics in the state space [5]:

$$\dot{\mathbf{x}}(t) = \mathbf{A}\mathbf{x}(t),$$

where A – state transition matrix (STM), and the variables x, y, z correspond to angular frequencies $\omega_i, \omega_j, \omega_k$, the indices of which show the imaginary coordinate axes of 4D space, $\mathbf{x}(t)$ – time-varying quaternion in vector representation, $\dot{\mathbf{x}}(t)$ – time derivative of a vector $\mathbf{x}(t)$.

Equation (6) is a matrix differential equation in time and state space. Multiplying the input vector by the STM produces an output vector in the form of a time derivative of the input vector. For constant values of angular frequencies $\omega_i, \omega_j, \omega_k$, the STM has the simplest form:

(1)

(2)

(5)

(6)

$$\mathbf{A} = \omega_i \mathbf{I} + \omega_j \mathbf{J} + \omega_k \mathbf{K} = \begin{bmatrix} 0 & \omega_i & \omega_j & \omega_k \\ -\omega_i & 0 & -\omega_k & \omega_j \\ -\omega_j & \omega_k & 0 & -\omega_i \\ -\omega_k & -\omega_j & \omega_i & 0 \end{bmatrix}.$$
 (7)

In more complex cases, it is possible to consider various combinations of STM degrees. A^n , as well as other linear functional transformations.

The derivative of a quaternion function with respect to a quaternion is calculated as

$$\frac{df(q)}{dq} = \left(\frac{\partial p}{\partial s} + \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}\right) + i\left(-\frac{\partial p}{\partial x} + \frac{\partial u}{\partial s} - \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}\right) + \\ + j\left(-\frac{\partial p}{\partial y} + \frac{\partial u}{\partial z} + \frac{\partial v}{\partial s} - \frac{\partial w}{\partial x}\right) + k\left(-\frac{\partial p}{\partial z} - \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} + \frac{\partial w}{\partial s}\right).$$
(8)

The derivative of a quaternion function with respect to its conjugate quaternion will have the form:

$$\frac{df(q)}{d\overline{q}} = \left(\frac{\partial p}{\partial s} - \frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} - \frac{\partial w}{\partial z}\right) + i\left(\frac{\partial p}{\partial x} + \frac{\partial u}{\partial s} + \frac{\partial v}{\partial z} - \frac{\partial w}{\partial y}\right) + \\
+ j\left(\frac{\partial p}{\partial y} - \frac{\partial u}{\partial z} + \frac{\partial v}{\partial s} + \frac{\partial w}{\partial x}\right) + k\left(\frac{\partial p}{\partial z} + \frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} + \frac{\partial w}{\partial s}\right).$$
(9)

In matrix representation, we write the derivative (8) of function (2) using abbreviated notations of partial derivatives as

$$\frac{df(q)}{dq} = \left(\mathbf{E}\frac{\partial}{\partial s} + \mathbf{I}\frac{\partial}{\partial x} + \mathbf{J}\frac{\partial}{\partial y} + \mathbf{K}\frac{\partial}{\partial z}\right) \begin{bmatrix} p \\ u \\ v \\ w \end{bmatrix} = \begin{bmatrix} \partial_s & \partial_x & \partial_y & \partial_z \\ -\partial_x & \partial_s & -\partial_z & \partial_y \\ -\partial_y & \partial_z & \partial_s & -\partial_x \\ -\partial_z & -\partial_y & \partial_x & \partial_s \end{bmatrix} \begin{bmatrix} p \\ u \\ v \\ w \end{bmatrix}.$$
(10)

The derivative with respect to the conjugate quaternion (9) in matrix representation has the form:

$$\frac{df(q)}{d\overline{q}} = \left(\mathbf{E}\frac{\partial}{\partial s} - \mathbf{I}\frac{\partial}{\partial x} - \mathbf{J}\frac{\partial}{\partial y} - \mathbf{K}\frac{\partial}{\partial z}\right) \begin{bmatrix} p \\ u \\ v \\ w \end{bmatrix} = \begin{bmatrix} \partial_s & -\partial_x & -\partial_y & -\partial_z \\ \partial_x & \partial_s & \partial_z & -\partial_y \\ \partial_y & -\partial_z & \partial_s & \partial_x \\ \partial_z & \partial_y & -\partial_x & \partial_s \end{bmatrix} \begin{bmatrix} p \\ u \\ v \\ w \end{bmatrix}.$$
(11)

As can be seen from (10) and (11), the matrix of the derivative with respect to the conjugate quaternion is equal to the transposed matrix of the derivative with respect to the quaternion.

In vector analysis, the complex Hamiltonian operator is often used:

$$\nabla = i\frac{\partial}{\partial x} + j\frac{\partial}{\partial y} + k\frac{\partial}{\partial z}.$$
(12)

Let us write the operator (12) using the basis matrices (4) in the form:

$$\nabla = \left(\mathbf{I}\frac{\partial}{\partial x} + \mathbf{J}\frac{\partial}{\partial y} + \mathbf{K}\frac{\partial}{\partial z}\right) = \begin{bmatrix} 0 & \partial_x & \partial_y & \partial_z \\ -\partial_x & 0 & -\partial_z & \partial_y \\ -\partial_y & \partial_z & 0 & -\partial_x \\ -\partial_z & -\partial_y & \partial_x & 0 \end{bmatrix}.$$
(13)

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If function (2) is harmonic with component $\cos \omega t$ and $\sin \omega t$, then when calculating the time derivatives of complex functions using (13), angular frequencies ω_i , ω_j , ω_k appear in the form of multipliers of the STM elements (7) of the dynamic's equation (6). Note that the angular frequencies are constant and the change in rotation angles occurs in accordance with the change in time.

The solution to the homogeneous linear matrix differential equation (6) will be the exponential of STM (7). The fundamental matrix of a three-frequency quaternion has the form [5]:

$$\Phi(\omega_{i}, \omega_{j}, \omega_{k}, t) = e^{At} = e^{(\omega_{i}I + \omega_{j}J + \omega_{k}K)t} =
= s(\omega_{i}, \omega_{j}, \omega_{k}, t)E + x(\omega_{i}, \omega_{j}, \omega_{k}, t)I + y(\omega_{i}, \omega_{j}, \omega_{k}, t)J + z(\omega_{i}, \omega_{j}, \omega_{k}, t)K,$$
(14)
where $s(\omega_{i}, \omega_{j}, \omega_{k}, t) = \cos \omega_{i}t \cos \omega_{j}t \cos \omega_{k}t - \sin \omega_{i}t \sin \omega_{j}t \sin \omega_{k}t,$

$$x(\omega_{i}, \omega_{j}, \omega_{k}, t) = \sin \omega_{i}t \cos \omega_{j}t \cos \omega_{k}t + \cos \omega_{i}t \sin \omega_{j}t \sin \omega_{k}t,$$

$$y(\omega_{i}, \omega_{j}, \omega_{k}, t) = \cos \omega_{i}t \sin \omega_{j}t \cos \omega_{k}t - \sin \omega_{i}t \cos \omega_{j}t \sin \omega_{k}t,$$

$$z(\omega_{i}, \omega_{j}, \omega_{k}, t) = \cos \omega_{i}t \cos \omega_{j}t \sin \omega_{k}t + \sin \omega_{i}t \sin \omega_{j}t \cos \omega_{k}t.$$
(15)

The fundamental matrix (14), taking into account (15), also represents a quaternion in matrix notation.

It is known that the fundamental matrix for three reference frequencies ω_i , ω_j , ω_k is decomposed into 4 single-frequency matrices [5]:

$$\Phi(\omega_{i}, \omega_{j}, \omega_{k}, t) =$$

$$= \Phi(\Omega_{1}, \Omega_{2}, \Omega_{3}, \Omega_{4}, t) = \Phi_{1}(\Omega_{1}, t) + \Phi_{2}(\Omega_{2}, t) + \Phi_{3}(\Omega_{3}, t) + \Phi_{4}(\Omega_{4}, t), \qquad (16)$$
where
$$\Phi_{1}(\Omega_{1}, t) = \frac{1}{4} \Big[(\cos(\Omega_{1}t) + \sin(\Omega_{1}t)) \mathbf{E} + (-\cos(\Omega_{1}t) + \sin(\Omega_{1}t)) \mathbf{I} + (\cos(\Omega_{1}t) + \sin(\Omega_{1}t)) \mathbf{J} + (-\cos(\Omega_{1}t) + \sin(\Omega_{1}t)) \mathbf{K}), \qquad (17)$$

$$+ (\cos(\Omega_{1}t) + \sin(\Omega_{2}t)) \mathbf{E} + (\cos(\Omega_{2}t) + \sin(\Omega_{2}t)) \mathbf{I} + (-\cos(\Omega_{2}t) - \sin(\Omega_{2}t)) \mathbf{K} \Big], \qquad (17)$$

$$\Phi_{3}(\Omega_{3}, t) = \frac{1}{4} \Big[(\cos(\Omega_{3}t) - \sin(\Omega_{2}t)) \mathbf{E} + (\cos(\Omega_{2}t) - \sin(\Omega_{2}t)) \mathbf{K} \Big], \qquad (18)$$

$$\Phi_{4}(\Omega_{4}, t) = \frac{1}{4} \Big[(\cos(\Omega_{4}t) + \sin(\Omega_{4}t)) \mathbf{E} + (-\cos(\Omega_{4}t) - \sin(\Omega_{4}t)) \mathbf{K} \Big], \qquad (18)$$

Thus, the quaternion space has a dimension of 4D. In this case, three coordinates are imaginary and one is real. For harmonic functions sine and cosine with different angular frequencies, each imaginary coordinate has its own angular frequency and a connection between time (rotation frequency) and space is formed. The STM for a dynamic state-space model is a differential operator and corresponds to the Hamiltonian operator. For harmonic functions, spatial differentiation corresponds to multiplying the elements of the quaternion by the angular frequencies of the corresponding spatial coordinate axes. The three-frequency fundamental matrix is a solution to the differential equation of dynamics in the state space. The three-frequency fundamental matrix for three reference frequencies is orthogonal and is decomposed into the sum (16) single-frequency orthogonal matrices (17) of combination frequencies (18).

3. Technique of Obtaining Maxwell's Quaternion Equations 3.1 Quaternion Electrostatic Equations

In electrostatics, charge q is considered in 3D space and charges of various bodies are calculated using geometric transformations [1]. In this case, the three coordinates of 3D space represent the geometric coordinates of bodies. This representation does not take into account the dimensionality of the charge in the energy space and, accordingly, the connection between the coordinates of the charge with its energy coordinates and the spatial electromagnetic fields that form the charge.

Let us represent the charge q as a quaternion and consider it in 4D energy space as a vector \mathbf{q} . Each axis of space will have a charge value and will be denoted by real numbers s, x, y, z, as in the algebraic representation of a quaternion (1). Since it is impossible to represent 4 axes on a plane, we will represent the scalar s as a sphere with coordinates for its imaginary axes. The radius of the sphere will determine the magnitude of the scalar [5]. To transform one 4D quaternion vector into another, a 4D matrix must be used. To satisfy the law of

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conservation of energy, in our case, charge, the matrix must be orthogonal.

Let us consider the imaginary unit of a single-frequency quaternion as the transformation matrix:

$$\hat{\mathbf{I}} = \frac{1}{\sqrt{3}} (\mathbf{I} + \mathbf{J} + \mathbf{K}) = \frac{1}{\sqrt{3}} \begin{bmatrix} 0 & 1 & 1 & 1 \\ -1 & 0 & -1 & 1 \\ -1 & 1 & 0 & -1 \\ -1 & -1 & 1 & 0 \end{bmatrix}.$$
(19)

Figure 1 shows two charged particles of different sign and magnitude. Red indicates a positive charge, and blue indicates a negative charge. A positively charged particle is represented by a vector $\mathbf{q}_1 = \begin{bmatrix} 5 & -2 & -3 & 2 \end{bmatrix}^T$ with the value of the scalar part of the charge *s*=5 and with the values of the imaginary charge coordinates in the matrix representation *x*=-2, *y*=-3, *z*=2. We use this vector as the initial state of the dynamic system (6). The ratio of the scalar part to the imaginary parts in the initial state can be anything. This ratio defines the norm of the charge vector. When multiplying by matrix (19), we obtain a charge $\mathbf{\hat{i}q_1} = \begin{bmatrix} -1.7 & 0 & -5.2 & -3.5 \end{bmatrix}^T$ in which the ratio of the scalar part to the imaginary parts is determined by the constancy of the norm of the resulting vectors. As can be seen, the norm of the vector $\mathbf{q_1}$ is equal to the norm of the vector $\mathbf{\hat{iq_1}}$. Consequently, the charge $\mathbf{q_1}$ cannot move to any place in the charge space, but only to the one in which it maintains its norm. This transition can be called a *quantum jump* of a 4D electron. Similar transformations are shown in Figure 2 for the charge vector $\mathbf{q_2} = \begin{bmatrix} -3 & 2 & -1 & 3 \end{bmatrix}^T$.

As is known, multiplication by an imaginary unit corresponds to a vector rotation. In our case, the rotation by matrix (19) can be divided into three rotations in 4D space, carried out by the basis matrices **I**, **J**, **K**.



Figure 1: The Resulting Vector [-1.7 0 -5.2 -3.5]



Figure 2: The Resulting Vector [2.3 4 1.2 0]

Based on the dynamic's equation (6), an electron can make many such jumps with different intervals depending on the change in time and the values of angular frequencies. Since the jumps occur without changing the energy, the movement occurs along an orbit with a constant radius, the length of which is determined by the energy (norm) of the initial vector.

Figure 3 shows the orbit of a single-frequency quaternion for the initial state $\mathbf{q}_1 = \begin{bmatrix} 5 & -2 & -3 & 2 \end{bmatrix}^T$ corresponding to Figure 1. The initial state is depicted as a sphere with black surface lines. Figure 4 shows the orbit for the initial state $\mathbf{q}_2 = \begin{bmatrix} -3 & 2 & -1 & 3 \end{bmatrix}^T$ shown in Figure 2.



Figure 3: Rotation of a Quaternion \mathbf{q}_1 Along an Orbit



Figure 4: Rotation of a Quaternion q_2 Along an Orbit

If rotation occurs without loss of energy, then it can last infinitely long. To move to another orbit with a different radius, it is necessary to obtain or spend energy. Note that the orbit is formed without using the fact of attraction of some bodies to others (gravity). When using a three-frequency quaternion (14), the rotation orbits will be more complex than with a single-frequency one. Figures 5 and 6 show examples of rotation orbits of the electrons discussed above [5].



Figure 5: Rotation of a Three-Frequency Quaternion \mathbf{q}_1 Along an Orbit



Figure 6: Rotation of a Three-Frequency Quaternion \mathbf{q}_2 Along an Orbit

It should be noted that the considered idea is consistent with M. Planck's hypothesis about quantum radiation and absorption of energy of elementary particles and its dependence on frequency.

According to Gauss's law, knowing the electric field strength E or electric flux D, the electron charge density ρ_q at a given location in 3D space can be found as an integral over a closed elementary surface dS with differential volume dv [1,2]:

$$\oint_{S} \mathbf{D} \cdot \mathbf{d}S = \int_{\text{vol}} \rho_{q} \, \mathbf{d}v = \left(\frac{\partial \mathbf{D}_{x}}{\partial x} + \frac{\partial \mathbf{D}_{y}}{\partial y} + \frac{\partial \mathbf{D}_{z}}{\partial z}\right) \mathbf{d}v = \rho_{q}$$

In other words, the real part of the quaternion (the charge scalar) is determined by the imaginary parts of the quaternion (the electric field strengths).

Using the concept of divergence, we write the expression for the charge density as

$$\operatorname{div} \mathbf{D} = \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z}, \ \operatorname{div} \mathbf{D} = \rho_q$$
(20)

If we represent the electric flux as a function of the quaternion (2) and apply the conjugate Hamiltonian operator (12) to the expression for the quaternion function, then the divergence (20) can be represented as their scalar product:

$$\overline{\nabla} \bullet \mathbf{f}(q) = -\nabla \bullet \mathbf{f}(q) = -\nabla \bullet \mathbf{D} = -\left(\partial_x u + \partial_y v + \partial_z w\right),\tag{21}$$

where $f(q) = \begin{bmatrix} u & v & w \end{bmatrix}^T$ – a function of a quaternion, represented as a vector of a pure quaternion, which is a mathematical model of the electrical field strength E or the electrical flux D in 3D space.

It is known that $D = \varepsilon_0 E$, where ε_0 is the permittivity of free space (electric constant) [1]. As can be seen from (21) and in accordance with Gauss's law, the charge is determined by the vectors of electrical intensity. When representing the electric field as quaternions, the strength E and the electric flux D are three-dimensional vectors in 3D and represent the imaginary part of the quaternion. This indicates that the charge or charge density of electrons in this representation is formed by electromagnetic waves and has a corresponding spectrum. According to the terminology of classical electrodynamics, such charges form "eddy" currents. However, charges can be formed by summing charges, in the form of scalars, and form "electric currents". Consequently, charge, as a scalar quantity, together with vector values of electrical intensity, exists not in 3D, but at least in 4D space.

Conservation of charge energy during transformations is equivalent to conservation of the quaternion norm and is ensured by the Cauchy-Riemann conditions (CRC), which must be satisfied by hypercomplex numbers. CRC regulate the increments of the amplitudes of real and imaginary numbers so that the norm of the vectors of these numbers, i.e., the radius when they move, is preserved. CRC are obtained by making the derivative of a quaternion function with respect to its conjugate quaternion (11) equal to zero. Physically, this means that there is no dissipation of energy in the direction perpendicular to the direction of motion. For expression (11) to be equal to zero, each term of the sum of expressions on orthogonal coordinates must be equal to zero. From here we obtain the CRC for the increment of the scalar part depending on the increments of the imaginary parts along different coordinate axes:

$\frac{\partial p}{\partial p}$	<u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u><u></u></u>	$\frac{\partial v}{\partial v}$	$\Box \frac{\partial w}{\partial w}$	∂p _	ди	∂v	∂w	
∂s	∂x	$\int \frac{\partial y}{\partial y}$	∂z ,	$\frac{\partial x}{\partial x}$	∂s	∂z	$+\overline{\partial y}$,	
∂p	ди	∂v	∂w	∂p	ди	∂v	дw	(22)
$\overline{\partial y}^{-}$	$\frac{\partial z}{\partial z}$	∂s	$\overline{\partial x}$,	$\overline{\partial z}$	$-\frac{\partial y}{\partial y}$	$\overline{\partial x}$	∂s .	

The presence of a real part and three imaginary parts in a 4D electron shows that electrons are formed both by scalars and by magnetic and electric fields. The electron and its charge are associated with wave processes that manifest themselves in the form of harmonic functions in the exponential mapping of hypercomplex space (14). In other words, the electron vector in energy space has a spectrum.

Figures 7-10 show the spectra of a three-frequency quaternion in the form of rectangular pulses of a 4D charge

 $\mathbf{q}_1 = \begin{bmatrix} 5 & -2 & -3 & 2 \end{bmatrix}^T$, with the rectangular pulses shifted by half the duration to the right so that the beginning of the pulses coincides with zero [6]. The spectrum of the 1st vector element is shown in red, the 2nd in blue, the 3rd in green, and the 4th in brown. The turquoise color shows the modulus of the vector of spectra of rectangular pulses symmetrical to 0. As can be seen from the graphs, the

spectra of pulses shifted in time by half the pulse duration differ from the sin(x)/x function for the spectra of pulses symmetrical to 0.

The spectra in Figures 7-10 show how the magnitude of the scalar part of the electron (red lines) and the magnitude of the imaginary parts (blue, green and brown lines) change depending on the change in combination frequencies (18): Ω_1 , Ω_2 , Ω_3 Ω_4 . Note that Parseval's equality holds for quaternion spectra.



Figure 7: Spectra of Charge Vector Elements q_1 for Combination Frequency Ω_1



Figure 9: Spectra of Charge Vector Elements q_1 for Combination Frequency Ω_3



Figure 10: Spectra of Charge Vector Elements q_1 for Combination Frequency Ω_4

Thus, an electric charge, when using a quaternion-based hypercomplex model, is represented by a 4D vector in the quaternion energy space, in which the first element is a scalar and the other three are vectors on imaginary spatial coordinates i, j, k or in matrix representation on **I**, **J**, **K**. Gauss's law defines charge density through the divergence of electric flux and shows how electric fields form a charge in the form of a scalar.

The time evolution of the charge vector values is represented by a dynamic model in state space, where the STM corresponds to the complex Hamiltonian operator. It is shown that an electron can perform quantum jumps with conservation of energy or rotate infinitely in an orbit with a constant radius in 4D. When using a three-frequency quaternion, the rotation orbit becomes more complex. Using the Cauchy-Riemann conditions, the linearity of the model allows us to obtain conditions that correspond to the law of conservation of energy and the principle of superposition.

The solution to the state-space dynamics equation, which is a linear homogeneous differential equation with constant coefficients, will be an exponential raised to the STM power. Since the charge model is based on a hypercomplex signal, then according to Euler's formula, hypercomplex numbers are represented in the form of harmonic functions sine and cosine with angular frequencies changing over time in a coordinate system with one real axis corresponding to the magnitude of the charge and with three imaginary (spatial) axes corresponding to the electric flow.

It is shown that using the quaternion Fourier transform it is possible to establish a connection between the magnitude of the charge elements and the value of the amplitudes and phases of the harmonics that form them, i.e., to find the frequency spectra of the charge vector elements. According to Parseval's equality for the quaternion Fourier transform, the charge value of the 4D electron vector is equal to the energy value of the harmonic functions of their generators.

3.2 Quaternion Magnetostatic Equations

Faraday introduced the concept of magnetic flux and used lines in space to visualize the magnetic field. He called these lines *flux lines* or *lines of force*, which show the direction of magnetic field strength at each point in space. In an electric field, it is clear that the electric flow is formed by particles that have charges and are called electrons. By analogy, we can say that the magnetic flux consists of particles, which we will call *rotors*.

As is known, the electron has spin, which is a form of angular momentum and is its fundamental property, like charge. Therefore, the electron acts as a charge in electrical interaction and as a rotor in magnetic interaction. The circulating current (rotor) creates a magnetic field. Each rotor produces magnetic intensity. Just like the charge, the angular velocity of the rotor ω can be either positive or negative and have different values. Opposite rotors attract each other, while negative rotors repel each other. Therefore, by analogy, the angular velocity $\pm \omega$ can be called the "*magnetic charge*".

According to Ampere's hypothesis, a magnet consists of many elementary magnetic rotors oriented in one direction. The magnet must be considered in vector space, since its magnetism depends on its location relative to the coordinate axes in 3D. The magnetic field vector starts at the center of the rotor and is perpendicular to it. Rotors with different directions of rotation are attracted to each other, forming a scalar. In this case, the imaginary components are compensated and the poles of the connected magnets disappear, forming a large magnet also with two poles.

Thus, an electron, which has an electric charge, is a scalar particle and forms charged bodies whose intensity is the same in all directions; a rotor is a vector particle and forms magnetic bodies whose intensity is perpendicular to the radius of rotation of the circulating currents. An electron, due to duality, manifests itself as a charge during electrical interaction, and as a rotor during magnetic interaction. A magnetic field can arise when an electric field circulates. The magnetic field of a coil is made up of the magnetic fields of individual turns. Electric flux is the flow of electron charges. Magnetic flux is the flow of electron spins (rotors).

In the classical theory of electromagnetism, it is stated that the lines of magnetic flux are closed and do not end at magnetic charges, since they do not exist, therefore $\oint_{S} \mathbf{B} dS = 0$. However, there is magnetic intensity and magnetic flux. Magnetic intensity and magnetic flux indicate the presence of rotors.

To measure the magnetic field strength, consider the following experiment. We drill a hole in the middle of a relatively large magnet and place it on an axis so that it rotates freely with almost no friction or inertia. We use the magnetic needle of a small compass as a point magnet. We will move the small compass to different points on the plane. The large magnet and the compass needle will rotate so that the opposite poles of the magnet and the needle point towards each other. That is, the compass needle will be affected by the force of the magnetic field of a large magnet. According to the Biot-Savart-Laplace law, the magnitude of the force will be the same on a circle whose center is located at the center of a large magnet.

This experiment is similar to the experiment with a large electric charge and with a point electric charge when measuring electrical intensity, which confirms Coulomb's law for charges. In the case of magnetic intensity, the Biot-Savart-Laplace law has a mathematical expression for magnetic intensity similar to Coulomb's law. Both show an inverse relationship with the square of the distance and a linear relationship between source and field. The difference between the experiment with charges and the experiment with magnets is

that the charge is a scalar quantity and its intensity manifests itself equally in all directions without rotation of the charge. The effect of magnetic force depends on the orientation of the magnet, so to maximize the effect of magnetic force, the magnet must be turned with the corresponding pole towards the location of the point magnet.

If we consider the electron spin as a source of magnetic tension, then the rotation of the spin occurs without mechanical movement and without inertia. Since the spins of the electrons that form the large magnet are always oriented toward the location of the point magnet, it is possible to use Gauss's law to determine the density of spins p_m or rotors in the magnet from the magnetic flux vector. This means that the total magnetic flux through an arbitrary closed surface with electrons creating a spin flux is equal to the closed circulating current, i.e., the rotor or spin:

$$\oint_{S} \mathbf{B} \cdot \mathbf{d} S = \int_{\text{vol}} \rho_{m} \, \mathbf{d} v = \left(\frac{\partial \mathbf{B}_{x}}{\partial x} + \frac{\partial \mathbf{B}_{y}}{\partial y} + \frac{\partial \mathbf{B}_{z}}{\partial z}\right) \mathbf{d} v = \rho_{m}.$$
(23)

Physically, this means that we measure the rotors inside the magnet without taking into account the external lines of force.

Using the concept of divergence, we write expression (23) for the density of rotors (spins) as

$$\operatorname{div} \mathbf{B} = \frac{\partial \mathbf{B}_x}{\partial x} + \frac{\partial \mathbf{B}_y}{\partial y} + \frac{\partial \mathbf{B}_z}{\partial z}, \ \operatorname{div} \mathbf{B} = \rho_m \tag{24}$$

If we apply the adjoint Hamiltonian operator (12) to the expression for the quaternion function in algebraic form (2), then (24) can be represented as their scalar product:

$$\overline{\nabla} \bullet \mathbf{f}(m) = -\nabla \bullet \mathbf{f}(m) = -\nabla \bullet \mathbf{B} = \partial_x u + \partial_y v + \partial_z w, \qquad (25)$$

where $f(q) = \begin{bmatrix} u & v & w \end{bmatrix}^T$ – a function of a quaternion, represented as a vector of a pure quaternion, which is a mathematical model of magnetic intensity or magnetic flux in 3D space.

It is known that the magnetic flux B is related to the magnetic intensity H as $B = \mu_0 H$, where μ_0 is the magnetic permeability of free space and $\sqrt{\mu_0 \varepsilon_0} = c$ is the speed of light.

In contrast to expression (21), the rotor is determined by the magnetic field strength H or magnetic flux B and shows how the magnetic flux forms the magnetic particle rotor (spin). It is known that when an electron moves in a magnetic field, it is affected by the Lorentz force, which makes it rotate. When modeling an electric field or a magnetic field using quaternions, we will have expressions (21) and (25) of the same form for the densities p_q and p_m . Accordingly, CRC (22) will also be valid for magnetic intensities. In addition, it will be possible to find spectra of 4D rotor vectors with different forms of amplitude changes over time and density in 4D space, similar to those presented in Figures 7–10.

3.3 Quaternion Electrodynamic Equations

The product of quaternions in algebraic representation is calculated as [4]

$$qp = (q_0 + iq_1 + jq_2 + kq_3)(p_0 + p_1i + p_2j + p_3k) = = (q_0p_0 - q_1p_1 - q_2p_2 - q_3p_3) + q_0(ip_1 + jp_2 + kp_3) + p_0(iq_1 + jq_2 + kq_3) + + i(q_2p_3 - q_3p_2) + j(q_3p_1 - q_1p_3) + k(q_1p_2 - q_2p_1).$$
(26)

Let us write pure quaternions in the form: $q = [0,q] = q_1i + q_2j + q_3k$ and $p = [0,p] = p_1i + p_2j + p_3k$. Using the notation made, we write the product (26) of the quaternions q and p as

$$qp = [q_0, q][p_0, p] = (q_0 p_0 - q \cdot p) + (q_0 p + p_0 q + q \times p),$$
(27)

The first term of the sum (26) $(q_0p_0 - (-q_1p_1 - q_2p_2 - q_3p_3)) = q_0p_0 - q \cdot p = q \cdot p$ is a real number representing the scalar product:

$$q \bullet p = (q_0 + q_1 i + q_2 j + q_3 k) \bullet (p_0 + p_1 i + p_2 j + p_3 k) = q_0 p_0 + q_1 p_1 + q_2 p_2 + q_3 p_3.$$

Let us write the second and third terms of the sum (26) as (26) as $q_0(ip_1 + jp_2 + kp_3) = q_0p$ and $p_0(iq_1 + jq_2 + kq_3) = p_0q$.

The fourth, fifth and sixth terms of the sum (26) correspond to the vector product of pure quaternions, i.e., circulations along the coordinate axes i, j, k:

$$i(q_2p_3-q_3p_2)+j(q_3p_1-q_1p_3)+k(q_1p_2-q_2p_1)=q\times p.$$

Let us represent expression (27) in vector form as a product $\mathbf{Q}^{\mathsf{T}}\mathbf{p}$, where matrix \mathbf{Q} is the matrix of the time derivative of the quaternion function with respect to the conjugate quaternion (11), changing over time, and the vector $\mathbf{p} = f(q) = \begin{bmatrix} p & u & v & w \end{bmatrix}^{\mathsf{T}}$ - quaternion function:

$$\mathbf{Q}^{\mathrm{T}}\mathbf{p} = \begin{bmatrix} \partial_{s,t}p - \partial_{x,t}u - \partial_{y,t}v - \partial_{z,t}w \\ \partial_{s,t}p + \partial_{s,t}u + \partial_{z,t}v - \partial_{y,t}w \\ \partial_{y,t}p - \partial_{z,t}u + \partial_{s,t}v + \partial_{x,t}w \\ \partial_{z,t}p + \partial_{y,t}u - \partial_{x,t}v + \partial_{s,t}w \end{bmatrix} \Leftrightarrow \begin{bmatrix} \frac{\mathrm{scalar}}{\partial_{s,t}p - \nabla \cdot \mathbf{f}(q)} \\ \partial_{s,t}p + \partial_{s,t}u - (\partial_{y,t}w - \partial_{z,t}v) \\ \partial_{y,t}p + \partial_{s,t}v - (\partial_{z,t}u - \partial_{x,t}w) \\ \partial_{s,t}p + \partial_{s,t}v - (\partial_{z,t}u - \partial_{x,t}w) \\ \partial_{s,t}p + \partial_{s,t}w - (\partial_{x,t}v - \partial_{y,t}u) \\ \nabla p \quad \overline{\partial_{s}\mathbf{f}(q)} \quad \nabla \times \mathbf{f}(q) \end{bmatrix},$$
(28)

where $\nabla = \begin{bmatrix} \partial_x & \partial_y & \partial_z \end{bmatrix}^T$, $\mathbf{f}(q) = \begin{bmatrix} u & v & w \end{bmatrix}$ – pure quaternions.

In expression (28), the notations for time derivatives retain the indices denoting the coordinate axes s, x, y, z in the matrix representation and showing that time derivatives are taken from complex functions. In the representations of functions from the quaternion p, u, v, w, time indices are absent and are implied by default. It is also necessary to remember that f(q) is a function of time.

In the resulting vector (28), instead of a function of a quaternion f(q) and, accordingly, a pure quaternion f(q), one can use the vectors of electric **E** or magnetic intensity **H** as quaternions or as pure quaternions E and H.

For the electrical intensity E or the electrical flux D, expression (28) is represented as

$$\begin{bmatrix} Gauss's law for E \\ \partial_{s,t} p_E - \nabla \cdot E \\ \rho_q \\ EMF of self & circulation \\ vector E \\ \partial_{x,t} p_E + \partial_{s,t} u_E & -(\partial_y w_E - \partial_z v_E) \\ \partial_{y,t} p_E + \partial_{s,t} v_E - (\partial_z u_E - \partial_x w_E) \\ \partial_{z,t} p_E + \partial_{s,t} w_E - (\partial_x v_E - \partial_y u_E) \\ \partial_{z,t} p_E + \partial_{s,t} w_E - (\partial_x v_E - \partial_y u_E) \\ \nabla_{p_E} & \partial_{s,t} E \\ \nabla_{x \in E} \end{bmatrix},$$
(29)

where $p_{\rm E}$, $u_{\rm E}$, $v_{\rm E}$, $w_{\rm E}$ – functions of electric field strength in quaternion representation, ${\rm E} = \begin{bmatrix} u_{\rm E} & v_{\rm E} & w_{\rm E} \end{bmatrix}^{\rm T}$ – pure quaternion of electric field vector **E**.

Based on CRC (22), which sets the requirements of the law of conservation of energy, each element of vector (29) must be equal to 0. Therefore, from (29) we obtain two equations:

$\nabla \cdot \mathbf{E} = \rho_q - \text{scalar equation},$	(30)
$\nabla \times \mathbf{E} = \partial_{s,t} \mathbf{E} + \nabla p_{\mathbf{E}}$. – vector equation.	(31)

Equation (30) represents Gauss's law (20) in the form of (21).

Faraday's law is known, written by Maxwell in differential form:

$$\nabla \times \mathbf{E} = -\partial_t \mathbf{B} \,. \tag{32}$$

Law (32) is usually explained as follows: the change in magnetic field B over time creates an electromotive force (EMF). EMF is considered induced because magnetic flux is involved in its creation. EMF is simply the voltage that produces current in a closed circuit. The minus sign in (32) shows that the EMF is in such a direction that the created current, whose magnetic flux is added to the original flux, reduces the magnitude of the magnetic flux. The statement that the induced voltage acts to oppose the change in magnetic flux is known as Lenz's law.

In the resulting representation (31), instead of the time-varying magnetic $flux \partial_t \mathbf{B}$, EMF $\partial_{s,t} \mathbf{E}$ is written, which creates a current in a closed circuit. As can be seen from expression (31), EMF $\partial_{s,t} \mathbf{E}$ is proportional to circulation $\nabla \times \mathbf{E}$ and is subtracted from it to obtain zero energy increment. In expression (31), in contrast to Maxwell's equation, there is an additional term $\nabla p_{\mathbf{E}}$ – the vector of derivatives of the scalar part of the electrical field strength along the imaginary coordinate axes in matrix representation. This is explained by the fact that Maxwell's equations are written for 3D space and, when represented in quaternion form, do not take into account the scalar part. The increment of the scalar part of the field strength vector along the imaginary coordinates, together with the increment of the electric field strength vector in the quaternion representation so that their sum is equal to zero, and the energy of the electric field strength vector in the quaternion representation is preserved.

For magnetic intensity **H** or magnetic flux **B**, expression (28) takes the form:

$$\begin{bmatrix} \underbrace{\operatorname{Gauss's law for H}}_{\partial_{s,t} p_{\mathbf{H}}} - \nabla \cdot \mathbf{H} \\ \underbrace{\partial_{s,t} p_{\mathbf{H}}}_{\rho_{m}} - \nabla \cdot \mathbf{H} \\ \underbrace{\partial_{s,t} p_{\mathbf{H}}}_{\partial_{m}} + \partial_{s,t} u_{\mathbf{H}} - \underbrace{\partial_{y,t} w_{\mathbf{H}} - \partial_{z,t} v_{\mathbf{H}}}_{\partial_{z,t} p_{\mathbf{H}}} \\ \underbrace{\partial_{y,t} p_{\mathbf{H}}}_{\delta_{y,t} p_{\mathbf{H}}} + \partial_{s,t} v_{\mathbf{H}} - \underbrace{\partial_{z,t} u_{\mathbf{H}} - \partial_{z,t} w_{\mathbf{H}}}_{\nabla \times \mathbf{H}} \\ \underbrace{\partial_{s,t} p_{\mathbf{H}}}_{\nabla \rho_{\mathbf{H}}} + \underbrace{\partial_{s,t} w_{\mathbf{H}}}_{\partial_{s} \mathbf{H}} - \underbrace{\partial_{x,t} w_{\mathbf{H}} - \partial_{y,t} u_{\mathbf{H}}}_{\nabla \times \mathbf{H}} \end{bmatrix},$$
(33)

where $p_{\mathbf{H}}$, u_{H} , v_{H} , w_{H} – functions of magnetic intensity in quaternion representation, $\mathbf{H} = \begin{bmatrix} u_{\mathrm{H}} & v_{\mathrm{H}} & w_{\mathrm{H}} \end{bmatrix}^{\mathrm{T}}$ – pure quaternion of the magnetic intensity vector \mathbf{H} .

Based on the law of conservation of energy, we obtain the following equations from (33):

$\nabla \cdot \mathbf{H} = \rho_m - \text{scalar equation},$	(34)
$\nabla \times \mathbf{H} = \partial_{s,t} \mathbf{H} + \nabla p_{\mathbf{H}}$ – vector equation.	(35)

Expression (34) corresponds to Gauss's law (24) for magnetic intensity or magnetic flux, written in the form (25). Let us recall that this expression corresponds to the density of rotors (electron spins) that form the magnetic intensity. In Maxwell's equations, the scalar product is assumed to be Δ ·H=0. If this is true, then magnetic intensity and magnetic flux cannot exist at different points in space, since only its presence indicates the presence of circulating currents (rotors).

Maxwell's equation for the circulation of the magnetic intensity vector is known [1,2]:

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \,.$$

Ampere established that the circulation of magnetic field strength creates an electric current $\nabla \times \mathbf{H} = \mathbf{J}$. This electric current will be induced, since a magnetic field is involved in its formation. **H**. In (35), instead of the induced (eddy) current **J** in the conductor, its EMF $\partial_{s,t} \mathbf{H}$ generator is shown. In addition, Maxwell added electric current to Ampere's equation, which is obtained by changing the electric flow. In (35) it is shown mathematically that electric current is formed by changing the scalar part of the magnetic intensity vector.

Thus, expressions for Maxwell's equations for a quaternion are obtained using the Cauchy-Riemann conditions. Basically, the resulting equations are similar to Maxwell's equations except for the presence of additional terms associated with the scalar part of the quaternion.

4. Conclusion

The equations of electrodynamics are obtained mathematically using the Cauchy-Riemann conditions for the quaternion. The equations

represent two systems of equations for electric and magnetic intensity. The equations are identical in form and consist of an equation for the scalar part and an equation for the vector 3D part. The scalar part of the system of equations for electrical field strength corresponds to Gauss's law for charge. The scalar part of the system of equations for magnetic intensity also corresponds to Gauss's law, in which the electron spin is considered as a scalar. The representation of a quaternion in 4D space with three spatial frequencies is consistent with M. Planck's hypothesis about quantum radiation and absorption of energy of elementary particles and its dependence on frequency.

The equations obtained mathematically for the vector parts of the quaternion basically correspond to Maxwell's equations for the circulation of electric and magnetic intensities, obtained on the basis of the experiments of Faraday and Ampere. The difference is that instead of induced (eddy) currents formed by electrical and magnetic intensities, electromotive forces created by these intensities are used. In addition, the resulting equations contain currents created by the scalar parts of the quaternion. As is known, Maxwell artificially, without relying on physical experiments, added a current formed by the change in electric flow over time to the equation of circulation of magnetic intensity. This fact confirms the necessity of a quaternion model in 4D space with a scalar coordinate in electrodynamics. Since the equations of electrodynamics are obtained using linear functions of the quaternion, these functions are also the solution of these equations.

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